

Geometry and Control of Human Eye Movements

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• In robotics

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Anatomy of the Eye



Muscle pulleys





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Muscle pulleys





Saccades: are the fastest eye movements (velocities: 30 ~ 700⁰/s and lasting for about 40ms). Aim is to precisely redirect the gaze to the target to have a stabilized image on the retina (diameter of about a degree). Ex: reading, a sudden eccentric sound. Happens under open-loop control.



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- Vestibular Ocular Reflex (VOR): compensates for the movement of the head ensuring a clear image of the target on retina.
- Vergence movements: are the ones where the target moves along the gaze axis toward or away from the eye. The eye, which has the target moves along the gaze axis, remains stationary.





• Planer eye movements

- Three-dimensional eye movements : Geometry
- Eye as a simple mechanical control system
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Planer Eye Movements

- To simplify experiments and analysis.
- Study of planer eye movements has led to a remarkable understanding of one-dimensional movements, from the muscle mechanics to the underlying neural control system.
- A detailed biomechanical model was proposed by Martin & Schovanec (1997), (M-S model) and studied saccadic eye movements. Sugathadasa et al.(2000) further investigated smooth-pursuit tracking problem.

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- **Polpitiya & Ghosh (2002)** proposed **"Learning Curves"** for open-loop saccadic movement control using the M-S model.






















The equation of motion for the eye globe can be written as:

 $J_g \ddot{\theta} + B_g \dot{\theta} + K_g \theta = F_{t_1} - F_{t_2}$

where J_G , B_G , and K_G denote the globe inertia, globe viscosity, and globe elasticity. J_g , B_g , K_g are obtained as $J_g = \frac{J_G}{980r(180/\pi)}$ with r denoting the radius of the eye globe.

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Model can be written in the form $\dot{x} = f(x) + g_1(x)u_1 + g_2(x)u_2$ where u_1 and u_2 are the neural inputs, let the state vector be $x^T(t) = [\theta, \dot{\theta}, l_{m1}, \dot{l}_{m1}, l_{m2}, \dot{l}_{m2}, F_{t_1}, F_{t_2}, a_1, a_2]$.

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$$f(x) = \begin{bmatrix} x_2 \\ \frac{1}{Jg}(x_7 - x_8 - B_g x_2 - K_g x_1) \\ x_4 \\ \frac{980}{M}(x_7 - F_{act}(x_3, x_4, x_9) - F_{pe}(x_3) - B_{pm}(\frac{180}{\pi r})x_4) \\ x_6 \\ \frac{980}{M}(x_8 - F_{act}(x_5, x_6, x_{10}) - F_{pe}(x_5) - B_{pm}(\frac{180}{\pi r})x_6) \\ K_t(x_7) \left[-x_2 - (\frac{180}{\pi r})x_4 \right] \\ K_t(x_8) \left[x_2 - (\frac{180}{\pi r})x_6 \right] \\ -\frac{x_9}{\tau_1} \\ -\frac{x_{10}}{\tau_1} \end{bmatrix}$$

$$\begin{split} g_1(x) &= (1/\tau_1) \left[0, 0, 0, 0, 0, 0, 0, 0, 1, 0 \right] \\ g_2(x) &= (1/\tau_1) \left[0, 0, 0, 0, 0, 0, 0, 0, 0, 1 \right]. \end{split}$$

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M-S Model: Simulations



Figure 1: Neuronal inputs and the resulting activation signals to the agonist and antagonist (10^0 saccade)

M-S Model: Simulations



Figure 2: Simulation of 10^o Saccade and the corresponding forces in the tendons

Learning Curves



Figure 3: "Learning Curves": Cubic Hermite interpolant splines developed from horizontal saccadic eye movements originating from the primary position. The bottom two figures demonstrate how the 'T' value changes with the initial gaze position.

Learning Curves



Figure 4: "Learning Curves": Cubic Hermite interpolant splines developed from horizontal saccadic eye movements originating from the primary position. The bottom two figures demonstrate how the 'T' value changes with the initial gaze position.

Outline of the talk

- Anatomy of the eye \checkmark
- Planer eye movements \checkmark
- Three-dimensional eye movements : Geometry
- Eye as a simple mechanical control system
- Optimal control of the eye
- Conclusions and future directions

- SO(3), the space of 3 × 3 rotation matrices, is the obvious choice for the configuration space.

$$SO(3) = \{ \mathbf{R} : \mathbb{R}^3 \to \mathbb{R}^3 \mid (\mathbf{R}x, \mathbf{R}y)_{\mathbb{R}^3} = (x, y)_{\mathbb{R}^3}, \det \mathbf{R} = 1 \}$$
$$= \{ \mathbf{R} : \mathbb{R}^3 \to \mathbb{R}^3 \mid \mathbf{R}\mathbf{R}^T = \mathrm{Id}, \det \mathbf{R} = 1 \}$$



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Eye as a *mechanical system with holonomic constraints*. Configuration space becomes a two dimensional submanifold of **SO**(3).

⇒ "Listing Space (*List*)"



- Space of quaternions are denoted by **Q**.
- $a \in \mathbf{Q}$ can be written as $a_0 \overrightarrow{\mathbf{1}} + a_1 \overrightarrow{\mathbf{i}} + a_2 \overrightarrow{\mathbf{j}} + a_3 \overrightarrow{\mathbf{k}}$.

•
$$\mathbf{vec}(a) = a_1 \overrightarrow{\mathbf{i}} + a_2 \overrightarrow{\mathbf{j}} + a_3 \overrightarrow{\mathbf{k}}$$

•
$$\mathbf{scal}(a) = a_0 \mathbf{1}$$

- The vector $a_1 \overrightarrow{\mathbf{i}} + a_2 \overrightarrow{\mathbf{j}} + a_3 \overrightarrow{\mathbf{k}}$ will be identified with $(a_1, a_2, a_3) \in \mathbb{R}^3$ without any explicit mention of it.
- Quaternion product: $p.q = p_0q_0 \mathbf{p} \cdot \mathbf{q} + p_0\mathbf{q} + q_0\mathbf{p} + \mathbf{p} \times \mathbf{q}$.

Thus we have maps,

vec :
$$\mathbf{Q} \to \mathbb{R}^3$$
, $a \mapsto (a_1, a_2, a_3)$,

and

$$\mathbf{scal}: \mathbf{Q} \to \mathbb{R}, \ a \mapsto a_0.$$



Each $q \in S^3$ (space of unit quaternions) can be written as

$$q = \cos(\alpha/2)\vec{1} + \sin(\alpha/2)n_1\vec{i} + \sin(\alpha/2)n_2\vec{j} + \sin(\alpha/2)n_3\vec{k}$$

where, $\alpha \in [0, \pi]$ and (n_1, n_2, n_3) is a unit vector in \mathbb{R}^3 .

Define

$$\mathbf{rot}: \mathrm{S}^3 \to \mathbf{SO}(3)$$

as the standard map from S^3 into SO(3) which maps $cos(\alpha/2)\vec{1} + sin(\alpha/2)n_1\vec{i} + sin(\alpha/2)n_2\vec{j} + sin(\alpha/2)n_3\vec{k}$ to a rotation around the axis n by a counterclockwise angle α .

There are two explicit ways of describing this map. First,

$$\mathbf{rot}(q)(v_1, v_2, v_3) = \mathbf{vec}(q.(v_1\overrightarrow{\mathbf{i}} + v_2\overrightarrow{\mathbf{j}} + v_3\overrightarrow{\mathbf{k}}).q^{-1}) \ .$$

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.

Second,

$$\mathbf{rot}(\mathbf{q}) = \begin{bmatrix} q_0^2 + q_1^2 - q_2^2 - q_3^2 & 2(q_1q_2 - q_0q_3) & 2(q_1q_3 + q_0q_2) \\ 2(q_1q_2 + q_0q_3) & q_0^2 + q_2^2 - q_1^2 - q_3^2 & 2(q_2q_3 - q_0q_1) \\ 2(q_1q_3 - q_0q_2) & 2(q_2q_3 + q_0q_1) & q_0^2 + q_3^2 - q_1^2 - q_2^2 \end{bmatrix} \in \mathbf{SO}(3).$$

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 - Local coordinates on *List*.
 - Riemannian metric on *List*.
 - Levi-Civita connection on *List*.
 - Geodesics on *List*.
 - Curvature on *List*.
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Local coordinates on *List*

Let x_3 axis is aligned with the normal gaze direction, then **Listing's law** amounts to a statement that all eye rotations have quaternion representations $q \in S^3$ with $q_3 = 0$.

List is diffeomorphic to \mathbb{P}^2 (antipodal points identified).

Local coordinates on *List*



(Note: this fails when $\phi = 0$ or $\phi = 2\pi$ since in both cases the the corresponding rotation is identity regardless of the value of θ)

Riemannian metric on *List*

Let's calculate the Riemannian metric on *List* induced from SO(3).

$$SO(3) = \{ \mathbf{R} : \mathbb{R}^3 \to \mathbb{R}^3 \mid (\mathbf{R}x, \mathbf{R}y)_{\mathbb{R}^3} = (x, y)_{\mathbb{R}^3}, \det \mathbf{R} = 1 \}$$
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The body *angular velocity* is defined as

$$\mathbf{\Omega}(t) = \mathbf{R}^T(t) \dot{\mathbf{R}}(t).$$

 $\mathbf{\Omega}(t)$ is a skew-symmetric matrix.

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 $\mathbf{\Omega}(t)$ is a skew-symmetric matrix.

Since

$$\dot{\mathbf{R}}(t) = \mathbf{R}(t)\mathbf{\Omega}(t), \qquad \mathbf{\Omega}^{T}(t) = -\mathbf{\Omega}(t),$$

the tangent space

$$\mathsf{T}_{R}\mathsf{SO}(3) = \{\mathbf{R}\mathbf{\Omega} \mid \mathbf{\Omega}^{T} = -\mathbf{\Omega}\}, \qquad \mathbf{R} \in \mathsf{SO}(3).$$

Then the tangent space to SO(3) at the identity:

$$\mathsf{T}_{\mathrm{Id}}\mathsf{SO}(3) = \{ \mathbf{\Omega} : \mathbb{R}^3 \to \mathbb{R}^3 \mid \mathbf{\Omega}^T = -\mathbf{\Omega} \} = so(3)$$

Note that the space so(3) is the Lie algebra of the Lie group SO(3).

Assuming that the eye as a perfect sphere, and its moment of inertia as $I_{3\times3}$, the left invariant Riemannian metric on SO(3) given by,

$$\left< \mathbf{\Omega}(e_i), \mathbf{\Omega}(e_j) \right>_I = \delta_{i,j} ,$$

where,

$$\mathbf{\Omega}(e_k) = \begin{bmatrix} 0 & \delta_{3,k} & -\delta_{2,k} \\ -\delta_{3,k} & 0 & \delta_{1,k} \\ \delta_{2,k} & -\delta_{1,k} & 0 \end{bmatrix},$$

and $\{\delta_{l,m}\}$ denotes the Kronecker delta function.



Now $\vec{i}, \vec{j}, \vec{k}$ is an orthonormal basis of $T_{\vec{1}}S^3$, and recall that $\mathbf{rot} : S^3 \to SO(3)$, then

$$\operatorname{rot}\left(\left[\begin{array}{c} \cos(t/2) \\ \sin(t/2) \\ 0 \\ 0 \end{array}\right]\right) = \operatorname{e}^{\operatorname{t}\Omega(e_1)}, \quad \operatorname{rot}\left(\left[\begin{array}{c} \cos(t/2) \\ 0 \\ \sin(t/2) \\ 0 \end{array}\right]\right) = \operatorname{e}^{\operatorname{t}\Omega(e_2)}, \quad \operatorname{rot}\left(\left[\begin{array}{c} \cos(t/2) \\ 0 \\ 0 \\ \sin(t/2) \end{array}\right]\right) = \operatorname{e}^{\operatorname{t}\Omega(e_3)}.$$



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Notice that,

$$\frac{d}{dt}\Big|_{t=0} \begin{bmatrix} \cos(t/2) \\ \sin(t/2) \\ 0 \\ 0 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \frac{\overrightarrow{\mathbf{i}}}{2}, \qquad \qquad \frac{d}{dt}\Big|_{t=0} e^{t\mathbf{\Omega}(e_1)} = \mathbf{\Omega}(e_1).$$



Now $\vec{i}, \vec{j}, \vec{k}$ is an orthonormal basis of $T_{\vec{1}}S^3$, and recall that **rot** : $S^3 \rightarrow SO(3)$, then

Therefore,

$$\operatorname{rot}_{\overrightarrow{i}}(\overrightarrow{i}/2) = \Omega(e_1), \qquad \operatorname{rot}_{\overrightarrow{i}}(\overrightarrow{j}/2) = \Omega(e_2), \qquad \operatorname{rot}_{\overrightarrow{i}}(\overrightarrow{k}/2) = \Omega(e_3).$$

Hence $\{\operatorname{rot}_{\overrightarrow{i1}}, \operatorname{rot}_{\overrightarrow{i1}}, \operatorname{rot$





The Riemannian metric on *List* has the form

$$g = ds^2 = \sum_{ij=1}^n g_{ij} dx_i dx_j$$

where

$$g_{ij} = \left\langle \frac{\partial}{\partial x_i}, \frac{\partial}{\partial x_j} \right\rangle, \qquad (x_1, x_2) = (\theta, \phi).$$



Let $\rho : [0, \pi] \times [0, 2\pi] \rightarrow S^3$,

$$\rho(\theta, \phi) = \begin{bmatrix} \cos(\phi/2) \\ \cos(\theta)\sin(\phi/2) \\ \sin(\theta)\sin(\phi/2) \\ 0 \end{bmatrix}$$





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$$\rho(\theta,\phi) = \begin{bmatrix} \cos(\phi/2) \\ \cos(\theta)\sin(\phi/2) \\ \sin(\theta)\sin(\phi/2) \\ 0 \end{bmatrix}.$$

$$\begin{array}{c} List & \xrightarrow{\rho} & S^{3} \\ \downarrow & \downarrow \\ T_{(\theta,\phi)}List & \xrightarrow{\rho_{*}} & T_{\rho(\theta,\phi)}S^{3} \end{array}$$

Then the Jacobian

$$\mathcal{J}(\rho)(\theta,\phi) = \begin{pmatrix} \rho_{*(\theta,\phi)}(\frac{\partial}{\partial\theta}) & \rho_{*(\theta,\phi)}(\frac{\partial}{\partial\phi}) \end{pmatrix} = \begin{pmatrix} 0 & -\frac{1}{2}sin(\phi/2) \\ -sin(\theta)sin(\phi/2) & \frac{1}{2}cos(\theta)cos(\phi/2) \\ cos(\theta)sin(\phi/2) & \frac{1}{2}sin(\theta)cos(\phi/2) \\ 0 & 0 \end{pmatrix}$$



Let $\rho : [0, \pi] \times [0, 2\pi] \to S^3$, $\rho(\theta, \phi) = \begin{bmatrix} \cos(\phi/2) \\ \cos(\theta)\sin(\phi/2) \\ \sin(\theta)\sin(\phi/2) \\ 0 \end{bmatrix}.$ $List \longrightarrow S^3$ $\downarrow \qquad \downarrow \qquad \downarrow$ $T_{(\theta, \phi)}List \longrightarrow T_{\rho(\theta, \phi)}S^3$

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Also notice that

$$\rho(\theta,\phi).\vec{\mathbf{i}} = \begin{bmatrix} -\cos(\theta)\sin(\phi/2) \\ \cos(\phi/2) \\ 0 \\ -\sin(\theta)\sin(phi/2) \end{bmatrix}, \quad \rho(\theta,\phi).\vec{\mathbf{j}} = \begin{bmatrix} -\sin(\theta)\sin(\phi/2) \\ 0 \\ \cos(\phi/2) \\ \cos(\theta)\sin(\phi/2) \end{bmatrix}, \quad \rho(\theta,\phi).\vec{\mathbf{k}} = \begin{bmatrix} 0 \\ \sin(\theta)\sin(\phi/2) \\ -\cos(\theta)\sin(\phi/2) \\ \cos(\phi/2) \end{bmatrix}$$



For θ = 0, it is easily observed that,

$$\begin{split} \rho_{*(0,\phi)}(\frac{\partial}{\partial \theta}) &= \sin(\phi/2)\cos(\phi/2)\rho(0,\phi).\overrightarrow{\mathbf{j}} - \sin^2(\phi/2)\rho(0,\phi).\overrightarrow{\mathbf{k}},\\ \rho_{*(0,\phi)}(\frac{\partial}{\partial \phi}) &= \frac{1}{2}\rho(0,\phi).\overrightarrow{\mathbf{i}}. \end{split}$$
Riemannian metric on *List*, (cont'd.)



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$$\rho_{*(0,\phi)}(\frac{\partial}{\partial \phi}) = \frac{1}{2}\rho(0,\phi).\vec{\mathbf{i}}.$$

Therefore

$$g_{11} = \left\langle \frac{\partial}{\partial \theta}, \frac{\partial}{\partial \theta} \right\rangle = 4sin^2(\phi/2),$$

$$g_{12} = \left\langle \frac{\partial}{\partial \theta}, \frac{\partial}{\partial \phi} \right\rangle = 0,$$

$$g_{22} = \left\langle \frac{\partial}{\partial \phi}, \frac{\partial}{\partial \phi} \right\rangle = 1.$$

Riemannian metric on *List*, (cont'd.)



For $\theta = 0$, it is easily observed that,

$$\rho_{*(0,\phi)}(\frac{\partial}{\partial \theta}) = \sin(\phi/2)\cos(\phi/2)\rho(0,\phi).\vec{\mathbf{j}} - \sin^2(\phi/2)\rho(0,\phi).\vec{\mathbf{k}},$$

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Thus, the Riemannian metric on *List*

$$g = 4sin^2(\phi/2)d\theta^2 + d\phi^2.$$

Levi-Civita connection on *List*

Riemannian connection, $\nabla : \mathfrak{X}(M) \to \mathfrak{X}(M)$ of a Riemannian manifold M, is uniquely defined by the *Koszul* formula

$$2\langle \nabla_X Y, Z \rangle = \mathscr{L}_X \langle Y, Z \rangle + \mathscr{L}_Y \langle X, Z \rangle - \mathscr{L}_Z \langle X, Y \rangle$$
$$-\langle X, [Y, Z] \rangle - \langle Y, [X, Z] \rangle + \langle Z, [X, Y] \rangle$$

A Riemannian connection ∇ has the following properties:

$$\begin{split} \nabla_{fX+gY} &= f \nabla_X + g \nabla_Y, \\ \nabla_X (aY+bZ) &= a \nabla_X Y + b \nabla_X Z, \\ \nabla_X fY &= \mathcal{L}_X fY + f \nabla_X Y, \end{split}$$

$$\begin{split} \nabla_X Y - \nabla_Y X &= [X, Y], \\ \mathcal{L}_X \langle Y, Z \rangle &= \langle \nabla_X Y, Z \rangle + \langle Y, \nabla_X Z \rangle \end{split}$$

for $X, Y, Z \in \mathfrak{X}(M)$, $f, g \in \mathfrak{F}(M)$ and $a, b \in \mathbb{R}$.

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for $X, Y, Z \in \mathfrak{X}(\mathsf{M})$, $f, g \in \mathfrak{F}(\mathsf{M})$ and $a, b \in \mathbb{R}$.

Using the subscripted coordinates (y_1, y_2) to denote (θ, ϕ) and *Christoffel symbols* Γ_{ii}^k

$$\nabla_{\partial y_i/\partial y_j} = \Gamma_{ij}^k \partial/\partial y_k,$$

Christoffel symbols are given by

$$\Gamma_{ij}^{k} = \sum_{h=1}^{2} \frac{g^{ih}}{2} \left\{ \frac{\partial g_{hj}}{\partial y_{k}} + \frac{\partial g_{hk}}{\partial y_{j}} - \frac{\partial g_{jk}}{\partial y_{h}} \right\} \quad i, j, k = 1, 2$$

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Levi-Civita connection on *List* (cont'd.)



Now

$$(g_{ij}) = \begin{pmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{pmatrix} = \begin{pmatrix} 4\sin^2(\phi/2) & 0 \\ 0 & 1 \end{pmatrix}, \quad (\text{lower } g\text{-}ij'\text{s})$$
$$(g^{ij}) = \begin{pmatrix} g^{11} & g^{12} \\ g^{21} & g^{22} \end{pmatrix} = \begin{pmatrix} \frac{1}{4\sin^2(\phi/2)} & 0 \\ 0 & 1 \end{pmatrix}. \quad (\text{upper } g\text{-}ij'\text{s})$$

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Levi-Civita connection on *List* (cont'd.)



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Thus, we obtain expressions for Christoffel symbols,

$$\begin{split} \Gamma_{11}^{1} &= 0, & \Gamma_{11}^{2} &= -sin(\phi), \\ \Gamma_{12}^{1} &= \frac{1}{2tan(\phi/2)}, & \Gamma_{21}^{1} &= \frac{1}{2tan(\phi/2)}, \\ \Gamma_{12}^{2} &= 0, & \Gamma_{21}^{2} &= 0, \\ \Gamma_{22}^{1} &= 0, & \Gamma_{22}^{2} &= 0. \end{split}$$

A geodesic is a curve whose length is the shortest distance between two points. Christoffel symbols can be used to compute geodesics. Let $\sigma(t) = (\theta(t), \phi(t))$ be a geodesic on *List*. Then

 $\nabla_{\dot{\sigma}(t)}\dot{\sigma}(t)=0,$

where

* * *

$$\dot{\sigma}(t) = \left(\dot{\theta}\frac{\partial}{\partial\theta} + \dot{\phi}\frac{\partial}{\partial\phi}\right)$$

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$$\nabla_{\dot{\sigma}(t)}\dot{\sigma}(t) = \nabla_{\left(\dot{\theta}\frac{\partial}{\partial\theta} + \dot{\phi}\frac{\partial}{\partial\phi}\right)} \left(\dot{\theta}\frac{\partial}{\partial\theta} + \dot{\phi}\frac{\partial}{\partial\phi}\right)$$

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$$\nabla_{\dot{\sigma}(t)}\dot{\sigma}(t) = \dot{\theta}\nabla_{\frac{\partial}{\partial\theta}}\left(\dot{\theta}\frac{\partial}{\partial\theta}\right) + \dot{\theta}\nabla_{\frac{\partial}{\partial\theta}}\left(\dot{\phi}\frac{\partial}{\partial\phi}\right) + \dot{\phi}\nabla_{\frac{\partial}{\partial\phi}}\left(\dot{\theta}\frac{\partial}{\partial\theta}\right) + \dot{\phi}\nabla_{\frac{\partial}{\partial\phi}}\left(\dot{\phi}\frac{\partial}{\partial\phi}\right)$$

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$$\nabla_{\frac{\partial}{\partial \theta}} \frac{\partial}{\partial \theta} = \sum_{k=1}^{2} \Gamma_{11}^{k} \frac{\partial}{\partial y_{k}} = -\sin(\phi) \frac{\partial}{\partial \phi}$$

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Therefore, the equations of geodesics

$$\ddot{\theta} + \frac{1}{tan(\phi/2)} \dot{\theta} \dot{\phi} = 0,$$
$$\ddot{\phi} - sin\phi \dot{\theta}^2 = 0.$$



The **curvature** \mathcal{R} of a Riemannian manifold (M, g) is a correspondence that associates to every pair $X, Y \in \mathfrak{X}(M)$ a mapping $\mathcal{R}(X, Y) : \mathfrak{X}(M) \to \mathfrak{X}(M)$ given by

 $\mathcal{R}(X,Y)Z = \nabla_Y \nabla_X Z - \nabla_X \nabla_Y Z + \nabla_{[X,Y]} Z, \quad Z \in \mathfrak{X}(\mathsf{M}),$

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From the Christoffel symbols for the basis $\{\partial_{\theta}, \partial_{\phi}\}, \mathcal{R},$

$$\mathcal{R}(\partial_{\theta}, \partial_{\phi})\partial_{\theta} = \nabla_{\partial_{\theta}} \nabla_{\partial_{\phi}} \partial_{\theta} - \nabla_{\partial_{\phi}} \nabla_{\partial_{\theta}} \partial_{\theta}, \qquad \text{since } [\partial_{\theta}, \partial_{\phi}] = 0, (\text{Note: } \partial_{\theta} = \frac{\partial}{\partial \theta}).$$

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$$\Re(X,Y)Z = \nabla_Y \nabla_X Z - \nabla_X \nabla_Y Z + \nabla_{[X,Y]} Z, \quad Z \in \mathfrak{X}(\mathsf{M}),$$

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This evaluates to,
$$\Re(\partial_{\theta}, \partial_{\phi})\partial_{\theta} = -\cos(\phi/2)\partial_{\theta}$$
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In particular, the Gauss curvature is given by,

$$\begin{aligned} K(\theta,\phi) &= \left\langle \mathcal{R}(\partial_{\theta},\partial_{\phi})\partial_{\phi},\partial_{\theta} \right\rangle / \left\langle \partial_{\theta},\partial_{\theta} \right\rangle \\ &= 1/4 \end{aligned}$$

Outline of the talk

- Anatomy of the eye \checkmark
- Planer eye movements \checkmark
- Three-dimensional eye movements : Geometry \checkmark
- Eye as a simple mechanical control system
- Optimal control of the eye
- Conclusions and future directions

Eye as a simple mechanical control system

A **"simple mechanical control system"** (see Smale, 1970) consists the following:

- a configuration manifold Q,
- Riemannian metric *g* on **Q** that defines the kinetic energy function on the tangent bundle of **Q**,
- external forces as functions on the tangent bundle,
- any constraints on the system,
- control forces on the system as covector fields on the configuration manifold.

Eye as a simple mechanical control system



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For the eye movement system, *List* is the configuration manifold.

 $g = 4sin^2(\phi/2)d\theta^2 + d\phi^2$ is the Riemannian metric on **List**.

Let the Lagrangian of the system be

$$\begin{split} L(\theta, \phi, \dot{\theta}, \dot{\phi}) &= \text{Kinetic Energy} - \text{Potential Energy} \\ &= \frac{1}{2} \left\| \dot{\theta} \frac{\partial}{\partial \theta} + \dot{\phi} \frac{\partial}{\partial \phi} \right\|^2 - V(\theta, \phi) \\ &= \frac{1}{2} \left\langle \dot{\theta} \frac{\partial}{\partial \theta}, \dot{\phi} \frac{\partial}{\partial \phi} \right\rangle - V(\theta, \phi) \end{split}$$

Let the Lagrangian of the system be

 $L(\theta, \phi, \dot{\theta}, \dot{\phi}) =$ Kinetic Energy – Potential Energy

$$= \frac{1}{2} \left\| \dot{\theta} \frac{\partial}{\partial \theta} + \dot{\phi} \frac{\partial}{\partial \phi} \right\|^2 - V(\theta, \phi)$$
$$= \frac{1}{2} \left\langle \dot{\theta} \frac{\partial}{\partial \theta}, \dot{\phi} \frac{\partial}{\partial \phi} \right\rangle - V(\theta, \phi)$$

Recall that

$$g_{11} = \langle \partial_{\theta}, \partial_{\theta} \rangle = 4sin^{2}(\phi/2),$$

$$g_{12} = \langle \partial_{\theta}, \partial_{\phi} \rangle = 0,$$

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$$= \frac{1}{2} \left\langle \dot{\theta} \frac{\partial}{\partial \theta}, \dot{\phi} \frac{\partial}{\partial \phi} \right\rangle - V(\theta, \phi)$$

$$L(\theta,\phi,\dot{\theta},\dot{\phi}) = 2\dot{\theta}^2 \sin^2(\phi/2) + \frac{1}{2}\dot{\phi}^2 - V(\theta,\phi)$$

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* * *

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$$L(\theta,\phi,\dot{\theta},\dot{\phi}) = 2\dot{\theta}^2 \sin^2(\phi/2) + \frac{1}{2}\dot{\phi}^2 - V(\theta,\phi)$$

Euler-Lagrange equations:

$$\frac{d}{dt}\frac{\partial L}{\partial \dot{q}^i} - \frac{\partial L}{\partial q^i} = F_i, \quad i = 1, \dots, n.$$

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Euler-Lagrange equations:

$$\frac{d}{dt}\frac{\partial L}{\partial \dot{q}^i} - \frac{\partial L}{\partial q^i} = F_i, \quad i = 1, \dots, n.$$

Therefore the equations of motion:

$$\ddot{\theta} + \dot{\theta}\dot{\phi}\cot(\phi/2) + \frac{1}{4}\csc^2(\phi/2)\frac{\partial}{\partial\theta}V = \frac{1}{4}\csc^2(\phi/2)\tau_{\theta}$$
$$\ddot{\phi} - \dot{\theta}^2\sin(\phi) + \frac{\partial}{\partial\phi}V = \tau_{\phi}.$$

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Outline of the talk

- Anatomy of the eye \checkmark
- Planer eye movements \checkmark
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- Optimal control of the eye
- Conclusions and future directions



Case I: Generalized torques, τ_{θ} , τ_{ϕ}

Let $V(\theta, \phi) = sin^2(\phi/2)$.

Equations of motion:

$$\ddot{\theta} + \dot{\theta}\dot{\phi}\cot(\phi/2) = \frac{1}{4}\csc^2(\phi/2)\tau_{\theta}$$
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Equations of motion:

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$$\ddot{\phi} - \dot{\theta}^2\sin(\phi) + \frac{1}{2}\sin(\phi) = \tau_{\phi}.$$

Let $[z_1, z_2, z_3, z_4]' = [\theta, \dot{\theta}, \phi, \dot{\phi}]'$, then

$$\frac{d}{dt} \begin{bmatrix} z_1 \\ z_2 \\ z_3 \\ z_4 \end{bmatrix} = \begin{bmatrix} z_2 \\ -z_2 z_4 cot(z_3/2) \\ z_4 \\ z_2^2 sin(z_3) - \frac{1}{2} sin(z_3) \end{bmatrix} + \begin{bmatrix} 0 \\ \frac{1}{4} csc^2(z_3/2) \\ 0 \\ 0 \end{bmatrix} \tau_{\theta} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \tau_{\phi}$$



We wish to control the state $(\theta, \dot{\theta}, \phi, \dot{\phi})$ from $(\theta_0, 0, \phi_0, 0)$ to $(\theta_1, 0, \phi_1, 0)$ in *T* unit of time, while minimizing the control energy,

$$\int_0^T \left[(\tau_\theta(t))^2 + (\tau_\phi(t))^2 \right] \mathrm{d}t.$$



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$$\int_0^T \left[(\tau_\theta(t))^2 + (\tau_\phi(t))^2 \right] \mathrm{d}t.$$

Lagrangian:

$$L=\frac{1}{2}\left((\tau_\theta(t))^2+(\tau_\phi(t))^2\right),$$

and denote the costate by λ . Construct the Hamiltonian

$$\begin{aligned} \mathcal{H}(z,\lambda) &= \lambda.\dot{z} - L(z) \\ &= \lambda_1 z_2 - \lambda_2 z_2 z_4 \cot(z_3/2) + \lambda_3 z_4 + \lambda_4 z_2^2 \sin(z_3) - \frac{1}{2} \lambda_4 \sin(z_3) \\ &\frac{\lambda_2}{4 \sin^2(z_3/2)} \tau_{\theta} + c \lambda_4 \tau_{\phi} + \frac{1}{2} \left((\tau_{\theta}(t))^2 + (\tau_{\phi}(t))^2 \right) \end{aligned}$$



Hamilton's principle:

$$\dot{q}^i = \frac{\partial \mathcal{H}}{\partial p_i}, \qquad \dot{p}_i = -\frac{\partial \mathcal{H}}{\partial q^i},$$

where $p_i = \frac{\partial L}{\partial \dot{q}^i}$, $i = 1, \dots, n$.



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Hamiltonian system:

$$\frac{d}{dt} \begin{bmatrix} z_1 \\ z_2 \\ z_3 \\ z_4 \\ \lambda_1 \\ \lambda_2 \\ \lambda_3 \\ \lambda_4 \end{bmatrix} = \begin{bmatrix} z_1 \\ -z_2 z_4 \cot(z_3/2) + 1/4 \sin^2(z_3/2) \tau_{\theta}^* \\ z_4 \\ z_2^2 \sin(z_3) - \frac{1}{2} \sin(z_3) + \tau_{\phi}^* \\ 0 \\ -\lambda_1 + \lambda_2 z_4 \cot(z_3/2) - 2\lambda_4 z_2 \sin(z_3) \\ -\frac{1}{2} \lambda_2 z_2 z_4 \csc^2(z_3/2) - \lambda_4 z_2^2 \cos z_3 + \frac{1}{2} \lambda_4 \cos(z_3) + \frac{1}{2} \lambda_2 \cot(z_3) \csc^2(z_3) \tau_{\theta}^* \\ \lambda_2 z_2 \cot(z_3/2) - \lambda_3 \end{bmatrix}$$


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According to Pontryagin Maximum Principle (PMP), we can obtain:

$$\begin{aligned} \tau_{\theta} &= -\frac{\lambda_2}{4\sin^2(z_3/2)}, \\ \tau_{\phi} &= -\lambda_4. \end{aligned}$$



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Thus the system becomes

$$\begin{bmatrix} \dot{z}_{1} \\ \dot{z}_{2} \\ \dot{z}_{3} \\ \dot{z}_{4} \\ \dot{\lambda}_{1} \\ \dot{\lambda}_{2} \\ \dot{\lambda}_{3} \\ \dot{\lambda}_{4} \end{bmatrix} = \begin{bmatrix} z_{2} \\ -z_{2}z_{4}\cot(z_{3}/2) - \frac{\lambda_{2}}{16}\csc^{4}(z_{3}/2) \\ z_{4} \\ z_{2}^{2}\sin(z_{3}) - \frac{1}{2}\sin(z_{3}) - \lambda_{4} \\ 0 \\ -\lambda_{1} + \lambda_{2}z_{4}\cot(z_{3}/2) - 2\lambda_{4}z_{2}\sin(z_{3}) \\ (-\frac{1}{2}\lambda_{2}z_{2}z_{4}\csc^{2}(z_{3}/2) - 2\lambda_{4}z_{2}^{2}\cos(z_{3}) + \frac{1}{2}\lambda_{4}\cos(z_{3}/2) - \frac{\lambda_{2}^{2}}{16}\csc^{4}(z_{3}/2)\cot(z_{3}/2)) \\ \lambda_{2}z_{2}\cot(z_{3}/2) - \lambda_{3} \end{bmatrix}$$



 $\dot{ heta}, \dot{\phi}$ and $au_{ heta}, au_{\phi}$

Case II: Simplified muscles

Each musculotendon consist of a linear spring with spring constant k_i , a damper with damping constant b_i , and an active force F_i .

Projecting the torques to *List*

 $\theta \longrightarrow \theta + \delta \theta, \phi \longrightarrow \phi.$

Virtual work by the spring: $k_i(l_i - l_{i_0})\delta l = k_i(l_i - l_{i_0})\frac{\partial l_i}{\partial \theta}d\theta$.



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Then the torque with the active force F_i with $C_i = k_i(l_i - l_{i_0}) + b_i(\dot{\theta}\frac{\partial l_i}{\partial \theta} + \dot{\phi}\frac{\partial l_i}{\partial \phi})$:

$$\tau_{\theta} = \sum_{i=1}^{6} \left[F_i + C_i \right] \frac{\partial l_i}{\partial \theta} \qquad \qquad \tau_{\phi} = \sum_{i=1}^{6} \left[F_i + C_i \right] \frac{\partial l_i}{\partial \phi}$$



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Optimal path and muscle forces, from $(\pi/6, \pi/6)$ to $(\pi/10, \pi/10)$



Case III: Hill-type muscles



Hill-type musculotendon

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Optimal path from $(\pi/5, \pi/6)$ to $(\pi/10, \pi/10)$



Superior and inferior rectus muscle activities

Lengths of (Eye) Rotations

* * *

$$\ell(\sigma) = \int_{a}^{b} \left\| \dot{\theta} \frac{\partial}{\partial \theta} + \dot{\phi} \frac{\partial}{\partial \phi} \right\| dt$$
$$= \int_{a}^{b} \sqrt{\dot{\theta}^{2} g_{11} + 2\dot{\theta} \dot{\phi} g_{12} + \dot{\phi}^{2} g_{22}} dt$$
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From	То	distance (radians)		
(θ,ϕ)	(θ,ϕ)	<i>SO</i> (3)	Geodesic	Min. energy
			on List	on List
$\left(\frac{\pi}{4},\frac{\pi}{6}\right)$	$\left(\frac{\pi}{8},\frac{\pi}{8}\right)$	0.219	0.222	0.324
$(\frac{\pi}{4},\frac{\pi}{4})$	$\left(\frac{\pi}{8},\frac{\pi}{6}\right)$	0.359	0.368	0.368
$\left(\frac{\pi}{6},\frac{\pi}{10}\right)$	$(\frac{\pi}{8},\frac{\pi}{4})$	0.476	0.480	0.482

Outline of the talk

- Anatomy of the eye \checkmark
- Planer eye movements \checkmark
- Three-dimensional eye movements : Geometry \checkmark
- Eye as a simple mechanical control system \checkmark
- Optimal control of the eye \checkmark
- Conclusions and future directions

Summary:

• Learning curves for planer saccadic eye movements.



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Future directions:

For fast eye movments (saccades), a better approach would be minimizing the time instead of the control. But higher dimensionality of the control (six muscle activities), makes it a harder problem. Simpler problem would be to solve the minimum-time problem with the generalized torques τ_θ, τ_φ.

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